



Today we will focus on the article "Sur les zéros de la fonction Zeta de Riemann" by Godfrey Harold Hardy [2]. It's a brief article, but the main result is considered the first step made toward solving the Riemann Hypothesis:

Theorem 1. *There is an infinite number of $t \in \mathbb{R}$ such that:*

$$\zeta\left(\frac{1}{2} + it\right) = 0.$$

Proof. Begin by considering the formula:

$$e^{-y} = \frac{1}{2\pi i} \int_{k-i\infty}^{k+i\infty} \Gamma(u)y^{-u} du \quad (1)$$

true for $Re(y) > 0$ and $k > 0$.

This is often referred to as the **Cahen-Mellin Integral**, it is obtained using the Mellin transformation \mathcal{M} . We will go briefly through the steps necessary to arrive to the equality, the details can be found in [1].

Given a function $f(y)$ defined on the positive real axis $0 < y < \infty$, the **Mellin transformation** \mathcal{M} is the operation mapping f in the function F defined in the complex plane as:

$$\mathcal{M}[f; u] \equiv F(u) := \int_0^{\infty} f(y)y^{u-1} dy$$

$F(u)$ is called the **Mellin transform** of f .

By definition, for $f(y) = e^{-py}$, $p > 0$, we have:

$$\mathcal{M}[e^{-py}; u] = \int_0^{\infty} e^{-py} y^{u-1} dy = \int_0^{\infty} e^{-y} \left(\frac{y}{p}\right)^{u-1} \frac{dy}{p} = p^{-u} \int_0^{\infty} e^{-y} y^{u-1} dy = p^{-u} \Gamma(u)$$

for $Re(u) > 0$.

Using Fourier's inversion theorem, one can find a direct way to invert Mellin's transformation in:

$$f(y) = \frac{1}{2\pi i} \int_{k-i\infty}^{k+i\infty} F(u)y^{-u} du.$$

Therefore, using the fact that $F(u) = \mathcal{M}[e^{-py}; u] = p^{-u}\Gamma(u)$, we obtain:

$$e^{-py} = \frac{1}{2\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} p^{-u} \Gamma(u) du. \quad (2)$$

Remark 1. *This is a more general form of equation 1, it is never mentioned in the article but, as we will see, it is necessary to proceed with the proof.*

Substitute $p = n^2$, $n \neq 0$, in equation 2 to obtain:

$$\begin{aligned} e^{-n^2 y} &= \frac{1}{2\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} n^{-2u} \Gamma(u) du \\ &\Downarrow \\ 2e^{-n^2 y} &= \frac{1}{\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} n^{-2u} \Gamma(u) du \\ &\Downarrow \\ 1 + 2e^{-n^2 y} &= 1 + \frac{1}{\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} n^{-2u} \Gamma(u) du \\ &\Downarrow \\ 1 + 2 \sum_{n=1}^{\infty} e^{-n^2 y} &= 1 + \sum_{n=1}^{\infty} \frac{1}{\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} n^{-2u} \Gamma(u) du \\ &\Downarrow \\ 1 + 2 \sum_{n=1}^{\infty} e^{-n^2 y} &= 1 + \frac{1}{\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} \sum_{n=1}^{\infty} n^{-2u} \Gamma(u) du \\ &\Downarrow \\ 1 + 2 \sum_{n=1}^{\infty} e^{-n^2 y} &= 1 + \frac{1}{\pi i} \int_{k-i\infty}^{k+i\infty} y^{-u} \zeta(2u) \Gamma(u) du. \end{aligned} \quad (3)$$

Notice that we are allowed to change the order of integration and series due to absolute convergence.

The above formula is only true for $k > \frac{1}{2}$, as the function $\zeta(2s)$ has a pole at $s = \frac{1}{2}$, we can therefore choose to move the integration contour left, if we remember to add, using the residue Theorem, $2\pi i$ times the residue at $s = \frac{1}{2}$, which is $\frac{1}{2\pi i} \sqrt{\frac{\pi}{y}}$, that is to say:

$$1 + \frac{1}{\pi i} \int_{\frac{1}{2}-i\infty}^{\frac{1}{2}+i\infty} y^{-u} \zeta(2u) \Gamma(u) du = 1 + \sqrt{\frac{\pi}{y}} + \frac{1}{\pi i} \int_{\frac{1}{4}-i\infty}^{\frac{1}{4}+i\infty} y^{-u} \zeta(2u) \Gamma(u) du. \quad (4)$$

Remember now the definition of Riemann's ξ function:

$$\xi(s) := \frac{s}{2} (s-1) \pi^{-\frac{s}{2}} \zeta(s) \Gamma\left(\frac{s}{2}\right)$$

and of the Ξ function:

$$\Xi(t) := \xi\left(\frac{1}{2} + it\right).$$

Using these definitions:

$$\begin{aligned} \Xi(2t) &= \xi\left(\frac{1}{2} + 2it\right) = \xi\left(2\left(\frac{1}{4} + it\right)\right) = \left(\frac{1}{4} + it\right)\left(\frac{1}{2} + 2it - 1\right)\pi^{-\frac{1}{4}-it}\zeta\left(\frac{1}{2} + 2it\right)\Gamma\left(\frac{1}{4} + it\right) \\ &\Downarrow \\ \zeta\left(\frac{1}{2} + 2it\right)\Gamma\left(\frac{1}{4} + it\right) &= \frac{\Xi(2t)\pi^{\frac{1}{4}+it}}{\left(\frac{1}{4} + it\right)\left(\frac{1}{2} + 2it - 1\right)}. \end{aligned} \quad (5)$$

Now, changing variable to $u = \frac{1}{4} + it$, the integral in equation 4 becomes:

$$\frac{1}{\pi i} \int_{\frac{1}{4}-i\infty}^{\frac{1}{4}+i\infty} y^{-u} \zeta(2u) \Gamma(u) du = \frac{1}{\pi i} \int_{-\infty}^{+\infty} \left(\frac{1}{y}\right)^{\frac{1}{4}+it} \zeta\left(\frac{1}{2} + 2it\right) \Gamma\left(\frac{1}{4} + it\right) idt \quad (6)$$

combining with equation 5 we find:

$$\begin{aligned} &= \frac{1}{\pi} \int_{-\infty}^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\left(\frac{1}{4} + it\right)\left(-\frac{1}{2} + 2it\right)} dt = \frac{1}{\pi} \int_{-\infty}^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{-\frac{1}{2}\left(\frac{1}{4} + 4t^2\right)} dt \\ &= -\frac{2}{\pi} \int_{-\infty}^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt. \end{aligned} \quad (7)$$

Therefore, equation 3 can be written as:

$$1 + 2 \sum_{n=1}^{\infty} e^{-n^2 y} = 1 + \sqrt{\frac{\pi}{y}} - \frac{2}{\pi} \int_{-\infty}^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt. \quad (8)$$

Notice that Ξ is an even function:

$$\Xi(t) = \xi\left(\frac{1}{2} + it\right) = \xi\left(1 - \frac{1}{2} - it\right) = \xi\left(\frac{1}{2} - it\right) = \Xi(-t)$$

where we used the known property of the ξ function: $\xi(s) = \xi(1-s)$ a rigorous proof of this can be found [on our site](#).

Using this property, we can rewrite the integral in 8 as:

$$\begin{aligned} &\frac{2}{\pi} \int_{-\infty}^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt = \frac{2}{\pi} \int_0^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt + \frac{2}{\pi} \int_{-\infty}^0 \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt \\ &= \frac{2}{\pi} \int_0^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt + \frac{2}{\pi} \int_{+\infty}^0 \left(\frac{\pi}{y}\right)^{\frac{1}{4}-it} \frac{\Xi(-2t)}{\frac{1}{4} + 4(-t)^2} (-dt) \\ &= \frac{2}{\pi} \int_0^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt + \frac{2}{\pi} \int_0^{+\infty} \left(\frac{\pi}{y}\right)^{\frac{1}{4}-it} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} dt \\ &= \frac{2}{\pi} \int_0^{+\infty} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} \left(\left(\frac{\pi}{y}\right)^{\frac{1}{4}+it} + \left(\frac{\pi}{y}\right)^{\frac{1}{4}-it} \right) dt. \end{aligned} \quad (9)$$

Replacing y with $y = \pi e^{i\alpha}$, $-\frac{1}{2}\pi < \alpha < \frac{1}{2}\pi$, we have:

$$\begin{aligned}
1 + 2 \sum_{n=1}^{\infty} e^{-n^2 \pi e^{i\alpha}} &= 1 + \sqrt{\frac{\pi}{\pi e^{i\alpha}}} - \frac{2}{\pi} \int_0^{+\infty} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} \left(\left(\frac{\pi}{\pi e^{i\alpha}} \right)^{\frac{1}{4} + it} + \left(\frac{\pi}{\pi e^{i\alpha}} \right)^{\frac{1}{4} - it} \right) dt \\
&\Downarrow \\
1 + 2 \sum_{n=1}^{\infty} e^{-n^2 \pi e^{i\alpha}} &= 1 + e^{-\frac{1}{2}i\alpha} - \frac{2}{\pi} \int_0^{+\infty} \frac{\Xi(2t)}{\frac{1}{4} + 4t^2} \left(e^{-\frac{1}{4}i\alpha} (e^{\alpha t} + e^{-\alpha t}) \right) dt \\
&\Downarrow \\
1 + 2 \sum_{n=1}^{\infty} e^{-n^2 \pi e^{i\alpha}} &= 1 + e^{-\frac{1}{2}i\alpha} - \frac{2e^{-\frac{1}{4}i\alpha}}{\pi} \int_0^{+\infty} \frac{(e^{\alpha t} + e^{-\alpha t}) \Xi(2t)}{\frac{1}{4} + 4t^2} dt \\
&\Downarrow \\
1 + 2 \sum_{n=1}^{\infty} e^{-n^2 \pi e^{i\alpha}} - 1 - e^{-\frac{1}{2}i\alpha} &= -\frac{2e^{-\frac{1}{4}i\alpha}}{\pi} \int_0^{+\infty} \frac{(e^{\alpha t} + e^{-\alpha t}) \Xi(2t)}{\frac{1}{4} + 4t^2} dt.
\end{aligned}$$

Multiply both sides by $-\frac{\pi}{2}e^{\frac{1}{4}i\alpha}$ to obtain

$$\begin{aligned}
-\frac{\pi}{2}e^{\frac{1}{4}i\alpha} \left(1 + 2 \sum_{n=1}^{\infty} e^{-n^2 \pi e^{i\alpha}} \right) + \frac{\pi}{2}e^{\frac{1}{4}i\alpha} + \frac{\pi}{2}e^{\frac{1}{4}i\alpha - \frac{1}{2}i\alpha} &= \int_0^{+\infty} \frac{(e^{\alpha t} + e^{-\alpha t}) \Xi(2t)}{\frac{1}{4} + 4t^2} dt \\
&\Downarrow \\
-\frac{\pi}{2}e^{\frac{1}{4}i\alpha} \left(1 + 2 \sum_{n=1}^{\infty} e^{-n^2 \pi e^{i\alpha}} \right) + \frac{\pi}{2} \left(e^{\frac{1}{4}i\alpha} + e^{-\frac{1}{4}i\alpha} \right) &= \int_0^{+\infty} \frac{(e^{\alpha t} + e^{-\alpha t}) \Xi(2t)}{\frac{1}{4} + 4t^2} dt. \tag{10}
\end{aligned}$$

In conclusion, we have:

$$\int_0^{+\infty} \frac{(e^{\alpha t} + e^{-\alpha t}) \Xi(2t)}{\frac{1}{4} + 4t^2} dt = -\frac{\pi}{2}e^{\frac{1}{4}i\alpha} (1 + 2\psi(e^{i\alpha})) + \pi \cos\left(\frac{1}{4}\alpha\right) \tag{11}$$

where

$$\psi(x) = \sum_{n=1}^{+\infty} e^{-n^2 \pi x}.$$

Remark 2. In the original article Hardy writes $F(q) = 1 + 2 \sum_{n=1}^{\infty} q^{n^2}$, instead of $\psi(x)$, where $q = e^{-\pi e^{i\alpha}}$.

He also notices that $F(q) = \theta_3(0, \tau)$; the function θ_3 is defined as

$\theta_3(z, q) = \sum_{-\infty}^{+\infty} q^{n^2} e^{2\pi iz}$ and is an example of the so called **Jacobi Theta functions**; key functions in many topics of mathematics.

We are using a different notation to match the one used in [3], given that the following part of the demonstration comes mainly from this other source.

The key formula used by Hardy in the proof of Theorem 1 is obtained from equation 11 differentiating both sides $2p$ times with respect to α .

Note that, provided $\alpha < \frac{1}{2}\pi$, we can differentiate the above integral with respect to α any number of times since $\zeta\left(\frac{1}{2} + it\right) = \mathcal{O}(t^A)$, $\Xi(t) = \mathcal{O}(t^A e^{-\frac{1}{2}\pi t})$.

We therefore have:

$$\int_0^\infty \frac{t^{2p}(e^{\alpha t} + e^{-\alpha t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt = \frac{(-1)^p \pi}{4^{2p}} \cos\left(\frac{1}{4}\alpha\right) - \left(\frac{d}{d\alpha}\right)^{2p} \left[\frac{\pi}{2} e^{\frac{1}{4}i\alpha} (1 + 2\psi(e^{i\alpha})) \right]. \quad (12)$$

The next step in the proof is to consider $\alpha \rightarrow \frac{1}{2}\pi$ and prove that the last term of equation 12 tends to 0, for all possible choices of p . Hardy does not prove this in detail, the demonstration that follows is an adaptation of the one in [3]:

For $\alpha \rightarrow \frac{1}{2}\pi$, $e^{i\alpha} \rightarrow e^{i\frac{1}{2}\pi} = i$, therefore we have to focus on $\psi(x)$ when $x \rightarrow i$.

It is a known property of the function $\psi(x)$ that:

$$2\psi(x) + 1 = \frac{1}{\sqrt{x}} \left[2\psi\left(\frac{1}{x}\right) + 1 \right]. \quad (13)$$

This equation comes up often while studying Analytic Number Theory, you can find a demonstration [on our site](#); (mind that there the function is denoted $\omega(x)$).

Equation 13 implies:

$$\psi(x) = \frac{1}{\sqrt{x}} \psi\left(\frac{1}{x}\right) + \frac{1}{2} \frac{1}{\sqrt{x}} - \frac{1}{2} \quad (14)$$

therefore:

$$\psi(i+\delta) = \sum_{n=1}^{\infty} e^{-n^2\pi(i+\delta)} = \sum_{n=1}^{\infty} e^{-n^2\pi i} e^{-n^2\pi\delta} = \sum_{n=1}^{\infty} (-1)^{-n^2} e^{-n^2\pi\delta} = \sum_{n=1}^{\infty} (-1)^n e^{-n^2\pi\delta}$$

where in the last equality we used the fact that n^2 and n have the same parity.

We can write:

$$\begin{aligned} \sum_{n=1}^{\infty} (-1)^n e^{-n^2\pi\delta} &= \sum_{n \in \mathbb{Z}_{>0}, n \text{ even}}^{\infty} e^{-n^2\pi\delta} - \sum_{n \in \mathbb{Z}_{>0}, n \text{ odd}}^{\infty} e^{-n^2\pi\delta} \\ &= \sum_{n=1}^{\infty} e^{-(2n)^2\pi\delta} - \left(\sum_{n=1}^{\infty} e^{-n^2\pi\delta} - \sum_{n \in \mathbb{Z}_{>0}, n \text{ even}}^{\infty} e^{-n^2\pi\delta} \right) \\ &= \sum_{n=1}^{\infty} e^{-n^2\pi 4\delta} - \left(\psi(\delta) - \sum_{n=1}^{\infty} e^{-(2n)^2\pi\delta} \right) \\ &= \psi(4\delta) - \psi(\delta) + \sum_{n=1}^{\infty} e^{-n^2\pi 4\delta} \\ &= 2\psi(4\delta) - \psi(\delta). \end{aligned} \quad (15)$$

Using now equation 14 we have:

$$\begin{aligned} 2\psi(4\delta) - \psi(\delta) &= \frac{2}{\sqrt{4\delta}}\psi\left(\frac{1}{4\delta}\right) + \frac{1}{\sqrt{4\delta}} - 1 - \frac{1}{\sqrt{\delta}}\psi\left(\frac{1}{\delta}\right) - \frac{1}{2\sqrt{\delta}} + \frac{1}{2} \\ &= \frac{1}{\sqrt{\delta}}\psi\left(\frac{1}{4\delta}\right) - \frac{1}{\sqrt{\delta}}\psi\left(\frac{1}{\delta}\right) - \frac{1}{2}. \end{aligned} \quad (16)$$

From this last equality, one sees that $\psi(i + \delta) \rightarrow 0$ for $\delta \rightarrow 0$, or in other words that $\psi(x)$ and all its derivatives tend to 0 for $x \rightarrow i$ along any route in an angle $|\arg(x - i)| < \frac{1}{2}\pi$.

Therefore, for $\alpha \rightarrow \frac{1}{2}\pi$, equation 12 tends to $\frac{(-1)^p \pi}{4^{2p}} \cos\left(\frac{1}{8}\pi\right)$:

$$\int_0^\infty \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt = \frac{(-1)^p \pi}{4^{2p}} \cos\left(\frac{1}{8}\pi\right). \quad (17)$$

The last step of the proof is the key one, Hardy supposes that we can find a $T > 1$ such that the function $\Xi(t)$ never changes its sign once we go over this value and this assumption brings us to a contradiction.

The function $\Xi(t)$ therefore never stops going from positive to negative, which proves that $\Xi(t) = \zeta\left(\frac{1}{2} + it\right)$ has infinitely many zeros.

Suppose that, for every $t > T > 1$, $\Xi(t)$ maintains one sign, say positive.

Consider now equation 17 with p odd:

$$\int_0^\infty \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt = \int_0^T \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt + \int_T^\infty \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt$$

↓

$$\frac{(-1)^p \pi}{4^{2p}} \cos\left(\frac{1}{8}\pi\right) - \int_0^T \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt = \int_T^\infty \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt$$

p is odd, therefore $\frac{(-1)^p \pi}{4^{2p}} \cos\left(\frac{1}{8}\pi\right) < 0$ and we have:

$$\int_T^\infty \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt < - \int_0^T \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt < KT^{2p} \quad (18)$$

where $K = - \int_0^T \frac{(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt$ does not depend on p .

This gives us a contradiction: By our assumption, $\Xi(2t) > \delta$ for some positive δ for every $t \in [2T, 2T + 1]$, therefore:

$$\int_T^\infty \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt > \int_{2T}^{2T+1} \frac{t^{2p}(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})\Xi(2t)}{\frac{1}{4} + 4t^2} dt > \delta(2T)^{2p} K_1 \quad (19)$$

where this time $K_1 = \int_{2T}^{2T+1} \frac{(e^{\frac{1}{2}\pi t} + e^{-\frac{1}{2}\pi t})}{\frac{1}{4} + 4t^2} dt$ is positive and still does not depend on p .

Now using 18 and 19 we have:

$$\begin{aligned}
 KT^{2p} &> \delta(2T)^{2p} K_1 \\
 &\Downarrow \\
 K &> \delta 2^{2p} K_1
 \end{aligned} \tag{20}$$

but this should be true for any odd integer p , while, if $K < 0$ it is obviously false and for $K > 0$ is false for any $p > \frac{\log_2\left(\frac{K}{K_1\delta}\right)}{2}$. □



Thank you!

**We hope this lesson has been beneficial in studying
this interesting topic.
For more lessons or demonstrations, visit our website.**

References

- [1] Jacqueline Bertrand, Pierre Bertrand, and Jean-Philippe Ovarlez. “The mellin transform”. In: *The transforms and applications handbook* (1995).
- [2] Godfrey Harold Hardy. “Sur les zéros de la fonction $\zeta(s)$ de Riemann”. In: *CR Acad. Sci. Paris* 158.1914 (1914), p. 1012.
- [3] Edward Charles Titchmarsh and David Rodney Heath-Brown. *The theory of the Riemann zeta-function*. Oxford university press, 1986.